

# Appendix A

## Examples concerning ergodicity

### A.1 On the definition of ergodicity

This section is devoted to the construction of a  $*$ -dynamical system  $(\mathfrak{A}, \varphi, \tau)$  with the property that if  $\|\tau(A) - A\|_\varphi = 0$ , then  $\|A - \alpha\|_\varphi = 0$  for some  $\alpha \in \mathbb{C}$ , but for which the fixed points of the operator  $U$  defined in Proposition 2.3.3 in terms of some cyclic representation, form a vector subspace of  $\mathfrak{H}$  with dimension greater than one. This will prove the necessity of a sequence, rather than a single element, in Definition 2.3.2, in order for Proposition 2.3.3 to hold.

First some general considerations. Consider a dense vector subspace  $\mathfrak{G}$  of a Hilbert space  $\mathfrak{H}$ , and let  $\mathfrak{L}(\mathfrak{H})$  be the bounded linear operators  $\mathfrak{H} \rightarrow \mathfrak{H}$ . Set

$$\mathfrak{A} := \{A|_{\mathfrak{G}} : A \in \mathfrak{L}(\mathfrak{H}), A\mathfrak{G} \subset \mathfrak{G} \text{ and } A^*\mathfrak{G} \subset \mathfrak{G}\}$$

where  $A|_{\mathfrak{G}}$  denotes the restriction of  $A$  to  $\mathfrak{G}$ , then  $\mathfrak{A}$  is clearly a vector subspace of  $\mathfrak{L}(\mathfrak{G})$ . For any  $A \in \mathfrak{A}$ , denote by  $\overline{A}$  the (unique) bounded linear extension of  $A$  to  $\mathfrak{H}$ . Now define an involution on  $\mathfrak{A}$  by

$$A^* := \overline{A}^*|_{\mathfrak{G}}$$

for all  $A \in \mathfrak{A}$ , then it is easily verified that  $\mathfrak{A}$  becomes a unital  $*$ -algebra. (For  $A, B \in \mathfrak{A}$  it is clear that  $AB$  is a bounded linear operator  $\mathfrak{G} \rightarrow \mathfrak{G}$  which therefore has the extension  $\overline{A}\overline{B} \in \mathfrak{L}(\mathfrak{H})$  for which  $\overline{A}\overline{B}\mathfrak{G} \subset \mathfrak{G}$  and  $(\overline{A}\overline{B})^*\mathfrak{G} = \overline{B}^*\overline{A}^*\mathfrak{G} \subset \mathfrak{G}$  by the definition of  $\mathfrak{A}$ . Hence  $AB \in \mathfrak{A}$ , which means that  $\mathfrak{A}$  is a subalgebra of  $\mathfrak{L}(\mathfrak{G})$ . Also,  $(AB)^* = (\overline{A}\overline{B})^*|_{\mathfrak{G}} = (\overline{B}^*\overline{A}^*)|_{\mathfrak{G}} = \overline{B}^*(\overline{A}^*|_{\mathfrak{G}}) = \overline{B}^*A^* = B^*A^*$ . Similarly for the other defining properties of an involution.) Note that for  $A \in \mathfrak{A}$  and  $x, y \in \mathfrak{G}$  we have

$$\langle x, Ay \rangle = \langle x, \overline{A}y \rangle = \langle \overline{A}^*x, y \rangle = \langle A^*x, y \rangle.$$

For a given norm one vector  $\Omega \in \mathfrak{G}$  we define a state  $\varphi$  on  $\mathfrak{A}$  by

$$\varphi(A) = \langle \Omega, A\Omega \rangle.$$

Next we construct a cyclic representation of  $(\mathfrak{A}, \varphi)$ . Let

$$\pi : \mathfrak{A} \rightarrow L(\mathfrak{G}) : A \mapsto A$$

then clearly  $\pi$  is linear with  $\pi(1) = 1$  and  $\pi(AB) = \pi(A)\pi(B)$ . Note that for any  $x, y \in \mathfrak{G}$  we have  $(x \otimes y)^* = y \otimes x$ , hence  $(x \otimes y)\mathfrak{G} \subset \mathfrak{G}$  and  $(x \otimes y)^*\mathfrak{G} \subset \mathfrak{G}$ , so  $(x \otimes y)|_{\mathfrak{G}} \in \mathfrak{A}$ . Now,  $\pi((x \otimes \Omega)|_{\mathfrak{G}})\Omega = x \langle \Omega, \Omega \rangle = x$ , hence  $\pi(\mathfrak{A})\Omega = \mathfrak{G}$ . Furthermore,  $\langle \pi(A)\Omega, \pi(B)\Omega \rangle = \langle A\Omega, B\Omega \rangle = \langle \Omega, A^*B\Omega \rangle = \varphi(A^*B)$ . Thus  $(\mathfrak{G}, \pi, \Omega)$  is a cyclic representation of  $(\mathfrak{A}, \varphi)$ .

Suppose we have a unitary operator  $U : \mathfrak{H} \rightarrow \mathfrak{H}$  such that  $U\mathfrak{G} = \mathfrak{G}$  and  $U\Omega = \Omega$ . Then  $U^*\mathfrak{G} = U^{-1}\mathfrak{G} = \mathfrak{G}$ , so  $V := U|_{\mathfrak{G}} \in \mathfrak{A}$ , and  $V^* = U^*|_{\mathfrak{G}}$ . It follows that  $VAV^* \in \mathfrak{A}$  for all  $A \in \mathfrak{A}$ , hence we can define a linear function  $\tau : \mathfrak{A} \rightarrow \mathfrak{A}$  by

$$\tau(A) = VAV^*.$$

Clearly  $V^*V = 1 = VV^*$ , so  $\tau(1) = 1$  and  $\varphi(\tau(A)^*\tau(A)) = \varphi(VA^*AV^*) = \langle U^*\Omega, A^*AU^*\Omega \rangle = \varphi(A^*A)$ , since  $U^*\Omega = U^{-1}\Omega = \Omega$ . Therefore  $(\mathfrak{A}, \varphi, \tau)$  is a \*-dynamical system. Note that  $U|_{\mathfrak{G}}$  satisfies equation (3.1) of Section 2.3, namely  $U\pi(A)\Omega = UA\Omega = UAU^*\Omega = \tau(A)\Omega = \pi(\tau(A))\Omega$ , hence  $U$  is the operator which appears in Proposition 2.3.3.

Assume  $\{x \in \mathfrak{G} : Ux = x\} = \mathbb{C}\Omega$ . If  $\|\tau(A) - A\|_{\varphi} = 0$ , it then follows for  $x = \iota(A)$ , with  $\iota$  given by equation (2.1) of Section 2.2, that  $\|Ux - x\| = \|\iota(\tau(A) - A)\| = \|\tau(A) - A\|_{\varphi} = 0$ , so  $x = \alpha\Omega$  for some  $\alpha \in \mathbb{C}$ . Therefore  $\|A - \alpha\|_{\varphi} = \|\iota(A - \alpha)\| = \|x - \alpha\Omega\| = 0$ .

In other words, assuming that the fixed points of  $U$  in  $\mathfrak{G}$  form the one-dimensional subspace  $\mathbb{C}\Omega$ , it follows that  $\|\tau(A) - A\|_{\varphi} = 0$  implies that  $\|A - \alpha\|_{\varphi} = 0$  for some  $\alpha \in \mathbb{C}$ .

It remains to construct an example of a  $U$  with all the properties mentioned above, whose fixed point space in  $\mathfrak{H}$  has dimension greater than one. The following example was constructed by L. Zsidó:

Let  $\mathfrak{H}$  be a separable Hilbert space with an orthonormal basis of the form

$$\{\Omega, y\} \cup \{u_k : k \in \mathbb{Z}\}$$

(that is to say, this is a total orthonormal set in  $\mathfrak{H}$ ) and define the linear operator  $U : \mathfrak{H} \rightarrow \mathfrak{H}$  via bounded linear extension by

$$\begin{aligned} U\Omega &= \Omega, \\ Uy &= y, \\ Uu_k &= u_{k+1}, \quad k \in \mathbb{Z}. \end{aligned}$$

A.1. ON THE DEFINITION OF ERGODICITY

75

Clearly  $U$  is isometric, while  $U\mathfrak{H}$  is dense in  $\mathfrak{H}$ , hence  $U$  is surjective, since  $\mathfrak{H}$  is complete. Since  $U$  is a surjective isometry, it is unitary. Let  $\mathfrak{G}$  be the linear span of

$$\{\Omega\} \cup \{y + u_k : k \in \mathbb{Z}\}.$$

Then  $U\mathfrak{G} = \mathfrak{G}$ . Furthermore,  $\mathfrak{G}$  is dense in  $\mathfrak{H}$ . Indeed,

$$\left\| y - \frac{1}{n} \sum_{k=1}^n (y + u_k) \right\| = \frac{1}{n} \left\| \sum_{k=1}^n u_k \right\| = \frac{1}{\sqrt{n}} \rightarrow 0$$

implies that  $y \in \overline{\mathfrak{G}}$ , the closure of  $\mathfrak{G}$ , hence also

$$u_k = (y + u_k) - y \in \overline{\mathfrak{G}}$$

for  $k \in \mathbb{Z}$ .

Next we show that

$$\{x \in \mathfrak{G} : Ux = x\} = \mathbb{C}\Omega. \tag{1.1}$$

If  $\alpha\Omega + \sum_{k=-n}^n \beta_k(y + u_k) \in \mathfrak{G}$  is left fixed by  $U$ , then

$$\alpha\Omega + \sum_{k=-n}^n \beta_k y + \sum_{k=-n}^n \beta_k u_{k+1} = \alpha\Omega + \sum_{k=-n}^n \beta_k y + \sum_{k=-n}^n \beta_k u_k$$

and it follows that  $\beta_{-n} = 0$ , and that  $\beta_{k+1} = \beta_k$  for  $k = -n, \dots, n-1$ . Thus

$$\alpha\Omega + \sum_{k=-n}^n \beta_k (y + u_k) = \alpha\Omega$$

proving (1.1).

On the other hand,

$$\{x \in \mathfrak{H} : Ux = x\}$$

clearly contains the two-dimensional vector space spanned by  $\Omega$  and  $y$ .

## A.2 An example of an ergodic system

Here we give the proof that Example 2.5.7 is indeed ergodic. It is clear that  $\tau$  is linear and that  $\tau(1) = 1$ . Let

$$A = \begin{pmatrix} a_{11} & a_{12} \\ a_{21} & a_{22} \end{pmatrix}$$

and

$$B = \begin{pmatrix} b_{11} & b_{12} \\ b_{21} & b_{22} \end{pmatrix}$$

be complex matrices. Then

$$\tau(A)^* = \begin{pmatrix} \overline{a_{22}} & \overline{c_2 a_{21}} \\ \overline{c_1 a_{12}} & \overline{a_{11}} \end{pmatrix}$$

and

$$\tau(A)^* \tau(A) = \begin{pmatrix} |a_{22}|^2 + |c_2 a_{21}|^2 & \overline{a_{22}} c_1 a_{12} + \overline{c_2 a_{21}} a_{11} \\ \overline{c_1 a_{12}} a_{22} + \overline{a_{11}} c_2 a_{21} & |c_1 a_{12}|^2 + |a_{11}|^2 \end{pmatrix}$$

while

$$A^* = \begin{pmatrix} \overline{a_{11}} & \overline{a_{21}} \\ \overline{a_{12}} & \overline{a_{22}} \end{pmatrix}$$

and

$$A^* A = \begin{pmatrix} |a_{11}|^2 + |a_{21}|^2 & \overline{a_{11}} a_{12} + \overline{a_{21}} a_{22} \\ \overline{a_{12}} a_{11} + \overline{a_{22}} a_{21} & |a_{12}|^2 + |a_{22}|^2 \end{pmatrix}$$

so

$$\begin{aligned} \varphi(\tau(A)^* \tau(A)) &= \frac{1}{2} (|a_{22}|^2 + |c_2 a_{21}|^2 + |c_1 a_{12}|^2 + |a_{11}|^2) \\ &\leq \frac{1}{2} (|a_{22}|^2 + |a_{21}|^2 + |a_{12}|^2 + |a_{11}|^2) \\ &= \varphi(A^* A) \end{aligned}$$

for all  $A$ , meaning that  $(\mathfrak{A}, \varphi, \tau)$  is a  $*$ -dynamical system, if and only if  $|c_1| \leq 1$  and  $|c_2| \leq 1$ , which is what we will assume.

Next we prove that it is ergodic. For even  $k \geq 0$  we have

$$\tau^k(B) = \begin{pmatrix} b_{11} & c_1^k b_{12} \\ c_2^k b_{21} & b_{22} \end{pmatrix}$$



## A.2. AN EXAMPLE OF AN ERGODIC SYSTEM

77

and therefore

$$A\tau^k(B) = \begin{pmatrix} a_{11}b_{11} + a_{12}c_2^k b_{21} & a_{11}c_1^k b_{12} + a_{12}b_{22} \\ a_{21}b_{11} + a_{22}c_2^k b_{21} & a_{21}c_1^k b_{12} + a_{22}b_{22} \end{pmatrix}$$

which means

$$\varphi(A\tau^k(B)) = \frac{1}{2} (a_{11}b_{11} + a_{12}c_2^k b_{21} + a_{21}c_1^k b_{12} + a_{22}b_{22}).$$

For odd  $k > 0$  we then get

$$\varphi(A\tau^k(B)) = \frac{1}{2} (a_{11}b_{22} + a_{12}c_2^k b_{21} + a_{21}c_1^k b_{12} + a_{22}b_{11})$$

by switching  $b_{11}$  and  $b_{22}$ . For  $c \in \mathbb{C}$  it is clear that  $U : \mathbb{C} \rightarrow \mathbb{C} : x \mapsto cx$  is a linear operator with  $\|U\| \leq 1$  if and only if  $|c| \leq 1$ , and for  $c \neq 1$  the only fixed point of  $U$  is 0, in which case

$$\frac{1}{n} \sum_{k=0}^{n-1} c^k x = \frac{1}{n} \sum_{k=0}^{n-1} U^k x \longrightarrow 0$$

for all  $x \in \mathbb{C}$  as  $n \rightarrow \infty$ , by the Mean Ergodic Theorem 2.4.1. Hence, for  $c_1 \neq 1$  and  $c_2 \neq 1$  it follows that

$$\begin{aligned} & \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=0}^{n-1} \varphi(A\tau^k(B)) \\ &= \lim_{n \rightarrow \infty} \frac{1}{n} \left\{ \begin{array}{l} \frac{n}{2} \left[ \frac{1}{2}(a_{11}b_{11} + a_{22}b_{22}) + \frac{1}{2}(a_{11}b_{22} + a_{22}b_{11}) \right] \text{ for } n \text{ even} \\ \frac{n-1}{2} \left[ \frac{1}{2}(a_{11}b_{11} + a_{22}b_{22}) + \frac{1}{2}(a_{11}b_{22} + a_{22}b_{11}) \right] + \frac{1}{2}(a_{11}b_{11} + a_{22}b_{22}) \text{ for } n \text{ odd} \end{array} \right\} \\ &= \left( \frac{a_{11} + a_{22}}{2} \right) \left( \frac{b_{11} + b_{22}}{2} \right) \\ &= \varphi(A)\varphi(B) \end{aligned}$$

which means that  $(\mathfrak{A}, \varphi, \tau)$  is ergodic, by Proposition 2.5.6(ii).

On the other hand, if  $c_1 = 1$  and  $c_2 \neq 1$ , then we have by a similar calculation that

$$\lim_{n \rightarrow \infty} \frac{1}{n} \sum_{k=0}^{n-1} \varphi(A\tau^k(B)) = \varphi(A)\varphi(B) + \frac{a_{21}b_{12}}{2}.$$

Likewise for the other cases where either  $c_1$  or  $c_2$  or both are equal to 1. So  $(\mathfrak{A}, \varphi, \tau)$  is ergodic if and only if  $c_1 \neq 1$  and  $c_2 \neq 1$ .

**A.2.1 Remark.** It is easily seen that  $\tau$  is not a homomorphism, namely

$$\tau(AB) = \begin{pmatrix} a_{21}b_{12} + a_{22}b_{22} & c_1(a_{11}b_{12} + a_{12}b_{22}) \\ c_2(a_{21}b_{11} + a_{22}b_{21}) & a_{11}b_{11} + a_{12}b_{21} \end{pmatrix}$$

while

$$\tau(A)\tau(B) = \begin{pmatrix} a_{22}b_{22} + c_1c_2a_{12}b_{21} & c_1(a_{22}b_{12} + a_{12}b_{11}) \\ c_2(a_{21}b_{22} + a_{11}b_{21}) & c_1c_2a_{21}b_{12} + a_{11}b_{11} \end{pmatrix}.$$

In fact, unless  $c_1c_2 = 1$ , it follows that we don't even have  $\tau(A^2) = \tau(A)^2$  for all  $A$ . Nor, for that matter, do we have  $\tau(A^*) = \tau(A)^*$  for all  $A$ , unless  $c_2 = \overline{c_1}$ . This is opposed to the situation for a measure theoretic dynamical system as defined in Section 2.1, where  $\tau$  in equation (1.1) of that section is always a  $*$ -homomorphism. It therefore makes sense not to assume that  $\tau$  is a  $*$ -homomorphism in Definition 2.3.1, since we now have an example where it isn't. ■

**A.2.2 Remark.** We note that  $\varphi(\tau(A)) = \varphi(A)$ , i.e.  $\varphi$  is  $\tau$ -invariant, but this fact in itself does not imply that  $\varphi(\tau(A)^*\tau(A)) \leq \varphi(A^*A)$ , since  $\tau$  is not a  $*$ -homomorphism, by Remark A.2.1.

Furthermore,  $\varphi(AB) = \varphi(BA)$  for all  $A, B \in \mathfrak{A}$ , so  $\varphi$  is commutative (so to speak) even though  $\mathfrak{A}$  is not. Also, while  $\tau(AB) \neq \tau(BA)$  for some  $A, B \in \mathfrak{A}$ , we still have  $\varphi(\tau(AB)) = \varphi(AB) = \varphi(BA) = \varphi(\tau(BA))$ , so  $\tau$  is noncommutative (so to speak), but with respect to  $\varphi$  it is again commutative. We conclude that while  $\mathfrak{A}$  is noncommutative,  $(\mathfrak{A}, \varphi, \tau)$  is still in many respects commutative simply because  $\varphi(AB) = \varphi(BA)$  for all  $A$  and  $B$ . ■

**A.2.3 Question.** Is there an example of an ergodic  $*$ -dynamical system  $(\mathfrak{A}, \varphi, \tau)$  in which  $\varphi(AB) \neq \varphi(BA)$  for some  $A, B \in \mathfrak{A}$ ? ■

## Bibliography

- [AM] S. L. Adler and A. C. Millard (1996). “Generalized quantum dynamics as pre-quantum mechanics”, *Nuclear Physics B* **473**, 199-244. (Section 1.7.)
- [AA] V. I. Arnold and A. Avez (1968). *Ergodic problems of classical mechanics*, W. A. Benjamin. (Remark 3.2.8.)
- [Ba] R. Balian (1991). *From Microphysics to Macrophysics: Methods and Applications of Statistical Physics Vol 1*, Springer-Verlag. (Section 1.7.)
- [BvN] G. Birkhoff and J. von Neumann (1936). “The logic of quantum mechanics”, *Annals of Mathematics* **37**, 823-843. (Remark 1.1.5.)
- [BL] P. Bocchieri and A. Loinger (1957). “Quantum Recurrence Theorem”, *Physical Review* **107**, 337-338. (Section 3.1.)
- [Bo] D. Bohm (1951). *Quantum theory*, Prentice-Hall (republished by Dover, 1989). (Example 1.6.5.)
- [BR] O. Bratteli and D.W. Robinson (1987). *Operator Algebras and Quantum Statistical Mechanics 1*, Springer-Verlag, 2nd edition. (Sections 1.2, 1.8 and 2.6, the GNS-construction 2.2.2, and Remarks 1.7.4 and 2.2.3.)
- [Bu] J. Bub (1977). “Von Neumann’s projection postulate as a probability conditionalization rule in quantum mechanics”, *Journal of Philosophical Logic* **6**, 381-390. (Section 1.5 and Remark 1.6.7.)
- [CFS] C. M. Caves, C. A. Fuchs and R. Schack (2002). “Quantum probabilities as Bayesian probabilities”, *Physical Review A* **65**, 022305. (Section 1.6.)
- [CDL] C. Cohen-Tannoudji, B. Diu and F. Laloë (1977). *Quantum mechanics*, Hermann/John Wiley & Sons. (Examples 1.6.4 and 1.9.3.)
- [Co] A. Connes (1994). *Noncommutative Geometry*, Academic Press. (Remark 1.9.2.)



- [DFR] S. Doplicher, K. Fredenhagen and J. E. Roberts (1995). “The Quantum Structure of Spacetime at the Planck Scale and Quantum Fields”, *Communications in Mathematical Physics* **172**, 187-220. (Remark 1.6.9.)
- [D1] R. Duvenhage (1999). *Quantum statistical mechanics, KMS states and Tomita-Takesaki theory*, MSc dissertation, University of Pretoria. (Remark 1.7.4.)
- [D2] R. Duvenhage (2002). “Recurrence in Quantum Mechanics”, *International Journal of Theoretical Physics* **41**, 45-61, quant-ph/0202023. (Chapter 1 and Sections 2.7 and 3.1.)
- [D3] R. Duvenhage (preprint, 2002). “The nature of information in quantum mechanics”, quant-ph/0203070. (Chapter 1, in particular Section 1.6.)
- [DS] R. Duvenhage and A. Ströh (preprint, 2001). “Recurrence and ergodicity in unital  $*$ -algebras”. (Chapter 2.)
- [EPR] A. Einstein, B. Podolsky and N. Rosen (1935). “Can Quantum-Mechanical Description of Physical Reality Be Considered Complete?”, *Physical Review* **47**, 777-780. (Example 1.6.5.)
- [Fe] W. Feller (1968). *An Introduction to Probability Theory and Its Applications, Volume I*, John Wiley & Sons, 3rd edition. (Section 1.5.)
- [Fi1] D. Finkelstein (1969). “Space-Time Code”, *Physical Review* **184**, 1261-1271. (Thought experiment 1.6.8.)
- [Fi2] D.R. Finkelstein (1996, corrected 1997). *Quantum Relativity: A Synthesis of the Ideas of Einstein and Heisenberg*, Springer-Verlag. (Remark 1.6.10.)
- [Fu] C. A. Fuchs (2001). “Quantum Foundations in the Light of Quantum Information”, in *Proceedings of the NATO Advanced Research Workshop on Decoherence and its Implications in Quantum Computation and Information Transfer*, edited by A. Gonis, Plenum, quant-ph/0106166; revised version “Quantum Mechanics as Quantum Information (and only a little more)”, quant-ph/0205039. (Section 1.6.)
- [FuP] C. A. Fuchs and A. Peres (2000). “Quantum Theory Needs No ‘Interpretation’”, *Physics Today* **53**(3), 70-71; “Quantum Theory - Interpretation, Formulation, Inspiration: Fuchs and Peres Reply”, *Physics Today* **53**(9), 14, 90. (Section 1.6.)



- [Ha] R. Haag (1996). *Local Quantum Physics: Fields, Particles, Algebras*, Springer-Verlag, 2nd edition. (Sections 1.2 and 1.9, Example 1.6.2 and Remarks 1.6.10 and 1.9.2.)
- [HH] T. Hogg and B.A. Huberman (1982). "Recurrence Phenomena in Quantum Dynamics", *Physical Review Letters* **48**, 711-714. (Section 3.1.)
- [Ho] A. S. Holevo (1989). "Limit Theorems for Repeated Measurements and Continuous Measurement Processes", in *Quantum Probability and Applications IV*, edited by L. Accardi and W. von Waldenfels, Springer-Verlag, pp. 229-255. (Example 1.6.4.)
- [Hu] R.I.G. Hughes (1989). *The Structure and Interpretation of Quantum Mechanics*, Harvard University Press. (Section 1.2.)
- [I] C.J. Isham (1995). *Lectures on Quantum Theory: Mathematical and Structural Foundations*, Imperial College Press. (Section 1.6 and Examples 1.6.5 and 1.6.6.)
- [J] E.T. Jaynes (1968). "Prior Probabilities", *IEEE Transactions on Systems Science and Cybernetics*, **SSC-4**, 227-241; reprinted in E.T. Jaynes (1983), *Papers on probability, statistics and statistical physics* (edited by R.D. Rosenkrantz), D. Reidel Publishing Company. (Subsection 1.8.1.)
- [KR1] R.V. Kadison and J.R. Ringrose (1986). *Fundamentals of the Theory of Operator Algebras, Volume I*, Academic Press. (Section 2.6.)
- [KR2] R.V. Kadison and J.R. Ringrose (1986). *Fundamentals of the Theory of Operator Algebras, Volume II*, Academic Press. (Section 1.7, Subsection 1.8.2 and Proposition 3.1.1.)
- [Kh] A. I. Khinchin (1949). *Mathematical Foundations of Statistical Mechanics*, Dover. (Remark 3.2.8.)
- [Ko] B.O. Koopman (1931). "Hamiltonian systems and transformations in Hilbert space", *Proceedings of the National Academy of Sciences of the United States of America* **17**, 315-318. (Section 1.3.)
- [Kre] E. Kreyszig (1978). *Introductory Functional Analysis with Applications*, John Wiley & Sons. (Example 1.9.3.)
- [La] E. C. Lance (1995). *Hilbert  $C^*$ -modules: A toolkit for operator algebraists*, Cambridge University Press. (Open Problem 2.7.5.)

- [Lu] G. Lüders (1951). “Über die Zustandsänderung durch den Messprozess”, *Annalen der Physik* **8**, 323-328. (Section 1.2.)
- [Ma] J. Marsden (1973). *Application of global analysis in mathematical physics*, Carleton Mathematical Lecture Notes No. 3. (Remarks 1.4.1 and 1.6.9.)
- [MS] B. Misra and E.C.G. Sudarshan (1977). “The Zeno’s Paradox in Quantum Theory”, *Journal of Mathematical Physics* **18**, 756-763. (Example 1.6.4.)
- [Mu] G.J. Murphy (1990). *C\*-algebras and operator theory*, Academic Press. (Section 1.2, Remarks 1.7.4 and 1.9.2, Theorem 2.7.3 and Lemma 3.2.5.)
- [NSZ] C.P. Niculescu, A. Ströh and L. Zsidó (to appear). “Noncommutative extensions of classical and multiple recurrence theorems”, *Journal of Operator Theory*. (Section 2.5.)
- [OP] M. Ohya and D. Petz (1993). *Quantum Entropy and Its Use*, Springer-Verlag. (Section 1.5.)
- [Om] R. Omnès (1994). *The Interpretation of Quantum Mechanics*, Princeton University Press. (Remark 1.1.2.)
- [Perc] I.C. Percival (1961). “Almost Periodicity and the Quantal  $H$  Theorem”, *Journal of Mathematical Physics* **2**, 235-239. (Section 3.1.)
- [Pete] K. Petersen (1983). *Ergodic theory*, Cambridge University Press. (Sections 2.1, 2.4 and 2.6, and Remark 3.2.8.)
- [Petz] D. Petz (1988). “Conditional Expectation in Quantum Probability”, in *Quantum Probability and Applications III*, edited by L. Accardi and W. von Waldenfels, Springer-Verlag, pp. 251-260. (Section 1.5.)
- [Rud] W. Rudin (1987). *Real and complex analysis*, McGraw-Hill, 3rd edition. (Section 1.3, Proposition 1.7.1, Subsection 1.8.1 and Remark 3.2.8.)
- [Rue] D. Ruelle (1969). *Statistical Mechanics: Rigorous Results*, W.A. Benjamin. (Remarks 1.3.1, 3.2.4 and 3.2.8.)
- [Sc] J. Schwinger (2001). *Quantum Mechanics: Symbolism of Atomic Measurements*, Springer-Verlag. (Section 1.6.)
- [SLB] S. Seshadri, S. Lakshmibala and V. Balakrishnan (1999). “Quantum revivals, geometric phases and circle map recurrences”, *Physics Letters A* **256**, 15-19.
- [Sm] L. Smolin (2000). *Three Roads to Quantum Gravity: A new understanding of space, time and the universe*, Weidenfeld & Nicolson. (Remark 1.6.9.)



- [SZ] S. Strătilă and L. Zsidó (1979). *Lectures on von Neumann algebras*, Editura Academiei and Abacus Press. (Remark 1.4.2, Section 1.7, Lemma 3.2.5 and Theorem 3.2.7.)
- [St] R.F. Streater (2000). Classical and quantum probability, *Journal of Mathematical Physics* **41**, 3556-3603. (Remark 1.6.10.)
- [Su] A. Sudbery (1986). *Quantum mechanics and the particles of nature: An outline for mathematicians*, Cambridge University Press. (Examples 1.6.3 and 1.6.4.)
- [T] W. Thirring (1983). *Quantum mechanics of large systems*, A Course in Mathematical Physics 4, Springer-Verlag. (Remark 3.2.8.)
- [vN1] J. von Neumann (1932). *Mathematische Grundlagen der Quantenmechanik*, Springer. English translation *Mathematical Foundations of Quantum Mechanics* by R.T. Beyer, Princeton 1955. (Remark 1.1.5 and Example 1.6.2.)
- [vN2] J. van Neumann (1932). “Proof of the quasi-ergodic hypothesis”, *Proceedings of the National Academy of Sciences of the United States of America* **18**, 70-82. (Remark 1.1.5.)
- [Wa] P. Walters (1982). *An Introduction to Ergodic Theory*, Springer-Verlag. (Section 2.1.)

The remark at the end of each reference indicates where in this thesis (apart from the Introduction) the reference appears.